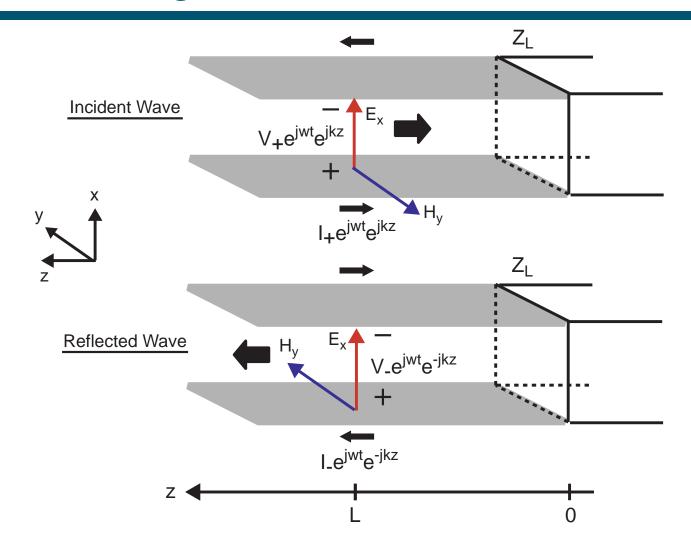
High Speed Communication Circuits and Systems Lecture 4 Generalized Reflection Coefficient, Smith Chart, Integrated Passive Components

Michael H. Perrott February 11, 2004

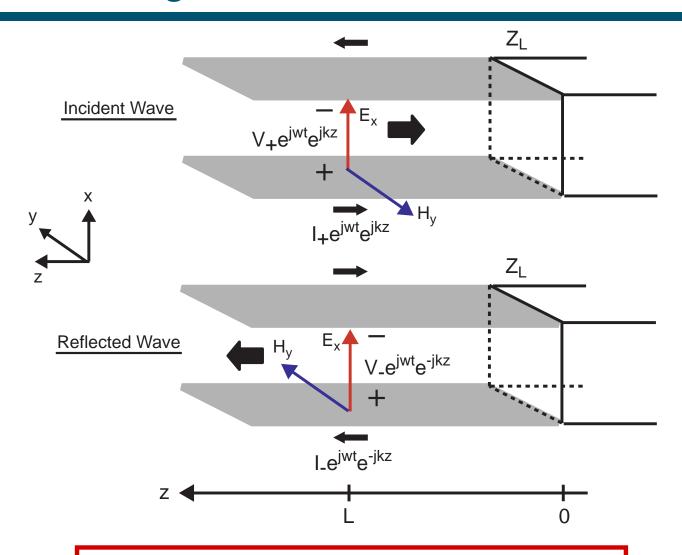
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Determine Voltage and Current At Different Positions



Incident and reflected waves must be added together

Determine Voltage and Current At Different Positions



$$V(z,t) = V_{+}e^{jwt}e^{jkz} + V_{-}e^{jwt}e^{-jkz}$$
$$I(z,t) = I_{+}e^{jwt}e^{jkz} - I_{-}e^{jwt}e^{-jkz}$$

Define Generalized Reflection Coefficient

$$V(z,t) = V_{+}e^{jwt}e^{jkz} + V_{-}e^{jwt}e^{-jkz}$$
$$I(z,t) = I_{+}e^{jwt}e^{jkz} - I_{-}e^{jwt}e^{-jkz}$$

$$V(z,t) = V_{+}e^{jwt}e^{jkz}\left(1 + \frac{V_{-}}{V_{+}}e^{-2jkz}\right)$$

$$V(z,t) = V_{+}e^{jwt}e^{jkz}\left(1 + \Gamma_{L}e^{-2jkz}\right)$$

$$V(z,t) = V_{+}e^{jwt}e^{jkz}\left(1 + \Gamma(z)\right)$$

Similarly:
$$I(z,t) = I_+ e^{jwt} e^{jkz} (1 - \Gamma(z))$$

$$\Rightarrow \Gamma(z) = \Gamma_L e^{-2jkz}$$

A Closer Look at $\Gamma(z)$

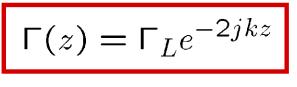
Recall Γ_L is

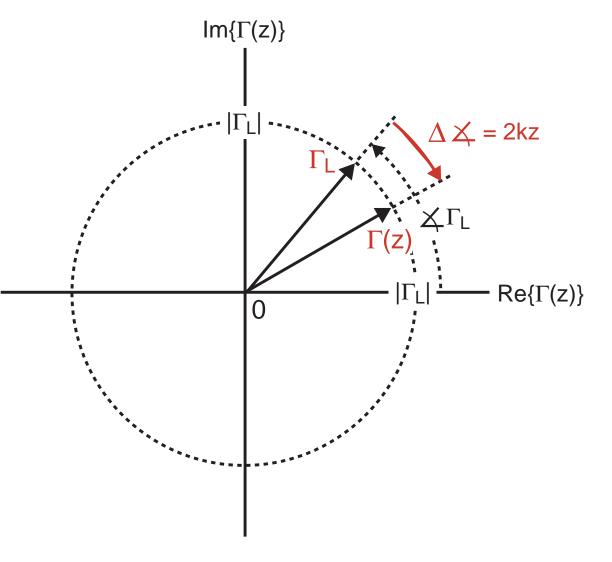
$$\Gamma_L = \frac{Z_L - Z_o}{Z_L + Z_o}$$

Note: $|\Gamma_L| \leq 1$

for
$$Re\{Z_L/Z_o\} \ge 0$$

We can view Γ(z) as a complex number that rotates clockwise as z (distance from the load) increases





Calculate $|V_{max}|$ and $|V_{min}|$ Across The Transmission Line

We found that

$$V(z,t) = V_{+}e^{jwt}e^{jkz}\left(1 + \Gamma(z)\right)$$

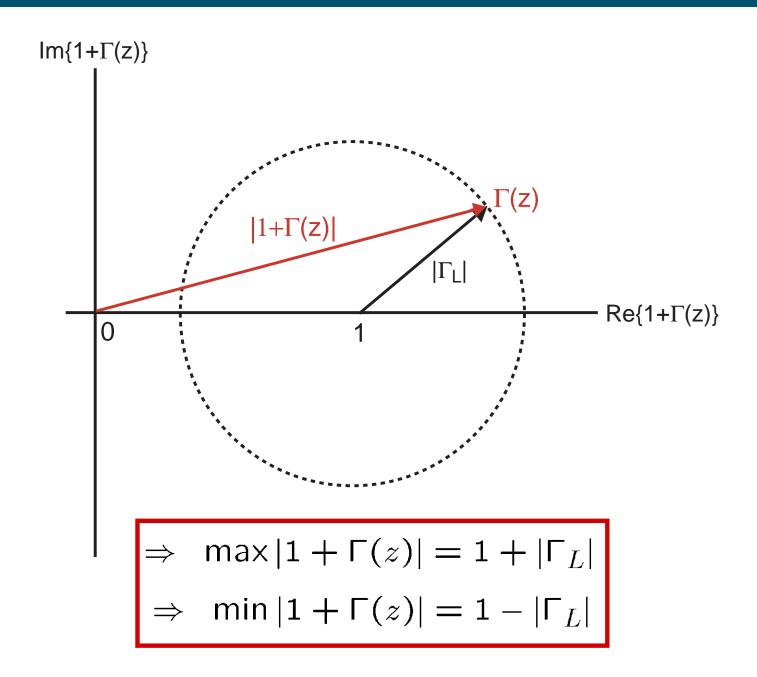
So that the max and min of V(z,t) are calculated as

$$\Rightarrow V_{max} = \max |V(z,t)| = |V_{+}| \max |1 + \Gamma(z)|$$

$$\Rightarrow V_{min} = \min |V(z,t)| = |V_{+}| \min |1 + \Gamma(z)|$$

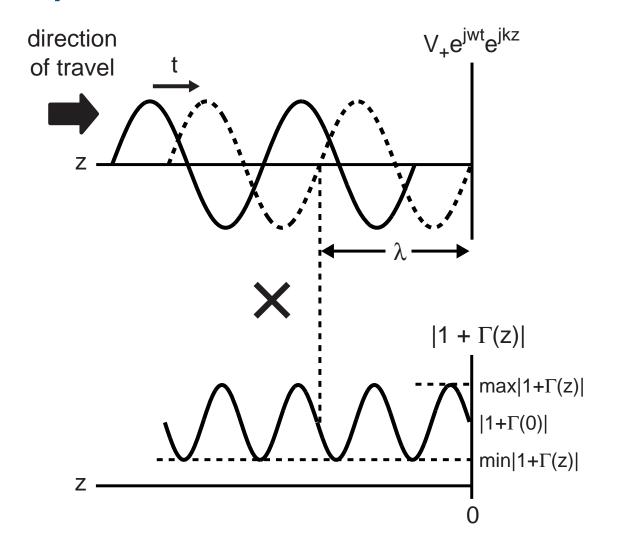
We can calculate this geometrically!

A Geometric View of $|1 + \Gamma(z)|$



Reflections Cause Amplitude to Vary Across Line

- Equation: $V(z,t) = V_{+}e^{jwt}e^{jkz}|1 + \Gamma(z)|e^{j\angle(1+\Gamma(z))}|$
- Graphical representation:



Voltage Standing Wave Ratio (VSWR)

Definition

VSWR =
$$\frac{V_{max}}{V_{min}} = \frac{|V_{+}|(1 + |\Gamma_{L}|)}{|V_{+}|(1 - |\Gamma_{L}|)} = \frac{1 + |\Gamma_{L}|}{1 - |\Gamma_{L}|}$$

For passive load (and line)

We can infer the magnitude of the reflection coefficient based on VSWR

$$|\Gamma_L| = rac{\mathsf{VSWR} - 1}{\mathsf{VSWR} + 1}$$

Reflections Influence Impedance Across The Line

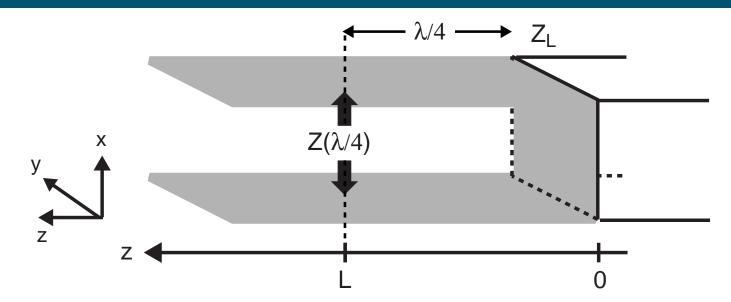
■ From Slide 4 $V(z,t) = V_{+}e^{jwt}e^{jkz}(1 + \Gamma(z))$ $I(z,t) = I_{+}e^{jwt}e^{jkz}(1 - \Gamma(z))$

$$\Rightarrow Z(z,t) = \frac{V_{+}(1+\Gamma(z))}{I_{+}(1-\Gamma(z))} = Z_{o}\frac{1+\Gamma(z)}{1-\Gamma(z)}$$

- Note: not a function of time! (only of distance from load)
- Alternatively $Z(z) = Z_o \frac{1 + \Gamma_L e^{-2jkz}}{1 \Gamma_L e^{-2jkz}}$
 - From Lecture 2: $\lambda = \frac{T}{\sqrt{\mu\epsilon}} = \frac{wT}{w\sqrt{\mu\epsilon}} = \frac{2\pi fT}{k} = \frac{2\pi}{k}$

$$Z(z) = Z_0 \frac{1 + \Gamma_L e^{-j(4\pi/\lambda)z}}{1 - \Gamma_L e^{-j(4\pi/\lambda)z}}$$

Example: $Z(\lambda/4)$ with Shorted Load



Calculate reflection coefficient

$$\Gamma_L = \frac{Z_L - Z_o}{Z_L + Z_o} = \frac{0 - Z_o}{0 + Z_o} = -1$$

Calculate generalized reflection coefficient

$$\Gamma(\lambda/4) = \Gamma_L e^{-j(4\pi/\lambda)(\lambda/4)} = \Gamma_L e^{-j\pi} = -\Gamma_L = 1$$

Calculate impedance $Z(\lambda/4) = Z_0 \frac{1 + \Gamma(z)}{1 - \Gamma(z)} = \infty$!

Generalize Relationship Between $Z(\lambda/4)$ and Z(0)

General formulation

$$Z(z) = Z_o \frac{1 + \Gamma_L e^{-j(4\pi/\lambda)z}}{1 - \Gamma_L e^{-j(4\pi/\lambda)z}}$$

At load (z=0)

$$Z_L = Z(0) = Z_0 \frac{1 + \Gamma_L}{1 - \Gamma_L}$$

• At quarter wavelength away ($z = \lambda/4$)

$$Z(\lambda/4) = Z_o \frac{1 - \Gamma_L}{1 + \Gamma_L} = \frac{Z_o^2}{Z_L}$$

- Impedance is inverted!
 - Shorts turn into opens
 - Capacitors turn into inductors

Now Look At Z(△) (Impedance Close to Load)

Impedance formula (∆ very small)

$$Z(\Delta) = Z_o \frac{1 + \Gamma_L e^{-2jk\Delta}}{1 - \Gamma_L e^{-2jk\Delta}}$$

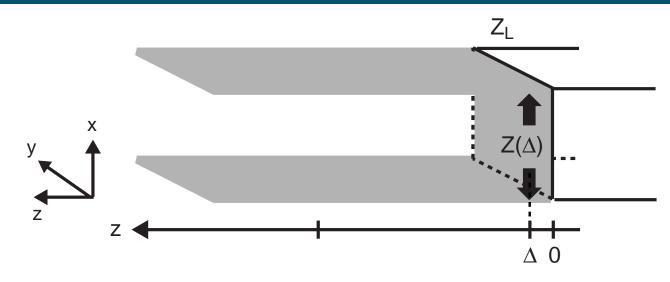
A useful approximation: $e^{-jx} \approx 1 - jx$ for $x \ll 1$

$$\Rightarrow e^{-2jk\Delta} \approx 1 - 2jk\Delta$$

- Recall from Lecture 2: $k = w\sqrt{LC}, \quad Z_o = \sqrt{\frac{L}{C}}$
- Overall approximation:

$$Z(\Delta) pprox \left(\sqrt{\frac{L}{C}}\right) rac{1 + \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}{1 - \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}$$

Example: Look At Z(\(\Delta\)) With Load Shorted

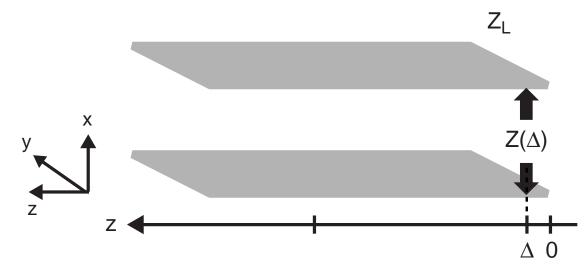


$$Z(\Delta) \approx \left(\sqrt{\frac{L}{C}}\right) \frac{1 + \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}{1 - \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}$$

- Reflection coefficient: $\Gamma_L = \frac{Z_L Z_o}{Z_L + Z_o} = \frac{0 Z_o}{0 + Z_o} = -1$
- Resulting impedance looks inductive!

$$Z(\Delta) pprox \left(\sqrt{\frac{L}{C}}\right) \frac{1 - (1 - 2jw\sqrt{LC}\Delta)}{1 + (1 - 2jw\sqrt{LC}\Delta)} pprox jwL\Delta$$

Example: Look At Z(△) With Load Open

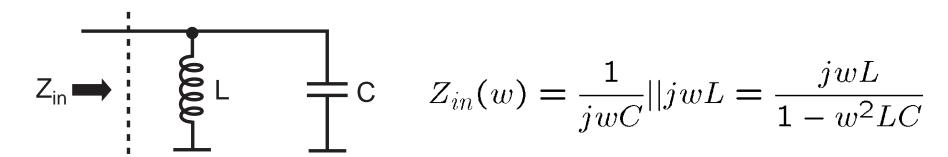


$$Z(\Delta) \approx \left(\sqrt{\frac{L}{C}}\right) \frac{1 + \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}{1 - \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}$$

- Reflection coefficient: $\Gamma_L = \frac{Z_L Z_o}{Z_L + Z_o} = \frac{\infty Z_o}{\infty + Z_o} = 1$
- Resulting impedance looks capacitive!

$$Z(\Delta) pprox \left(\sqrt{\frac{L}{C}}\right) rac{1 + (1 - 2jw\sqrt{LC}\Delta)}{1 - (1 - 2jw\sqrt{LC}\Delta)} pprox rac{1}{jwC\Delta}$$

Consider an Ideal LC Tank Circuit



Calculate input impedance about resonance

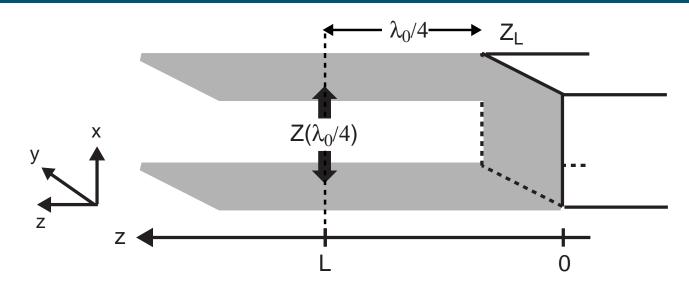
Consider
$$w = w_o + \Delta w$$
, where $w_o = \frac{1}{\sqrt{LC}}$

$$Z_{in}(\Delta w) = \frac{j(w_o + \Delta w)L}{1 - (w_o + \Delta w)^2 LC}$$

$$= \frac{j(w_o + \Delta w)L}{1 - w_o^2 LC - 2\Delta w(w_o LC) - \Delta w^2 LC}$$
= 0 negligible

$$\Rightarrow Z_{in}(\Delta w) \approx \frac{j(w_o + \Delta w)L}{-2\Delta w(w_o LC)} \approx \frac{jw_o L}{-2\Delta w(w_o LC)} = -\frac{j}{2}\sqrt{\frac{L}{C}}\left(\frac{w_o}{\Delta w}\right)$$

Transmission Line Version: $Z(\lambda_0/4)$ with Shorted Load



As previously calculated

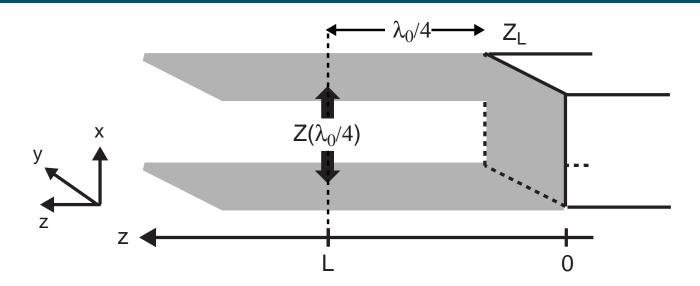
$$\Gamma_L = \frac{Z_L - Z_o}{Z_L + Z_o} = \frac{0 - Z_o}{0 + Z_o} = -1$$

Impedance calculation

$$Z(z) = Z_o \frac{1 + \Gamma(z)}{1 - \Gamma(z)}$$
, where $\Gamma(z) = \Gamma_L e^{-j(4\pi/\lambda)z}$

Relate λ to frequency $\lambda = \frac{1}{f\sqrt{\mu\epsilon}} = \frac{1}{(f_o + \Delta f)\sqrt{\mu\epsilon}}$

Calculate $Z(\Delta f)$ – Step 1



Wavelength as a function of ∆ f

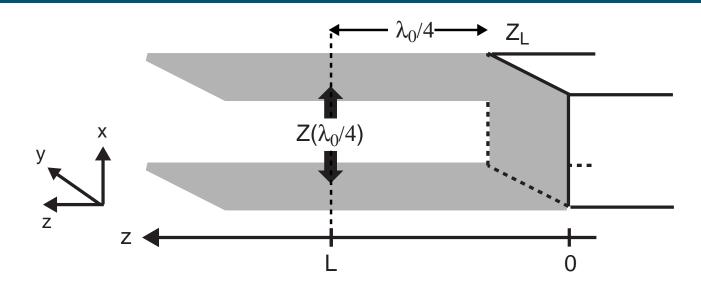
$$\lambda = \frac{1}{(f_o + \Delta f)\sqrt{\mu\epsilon}} = \frac{1}{f_o\sqrt{\mu\epsilon}(1 + \Delta f/f_o)} = \frac{\lambda_o}{1 + \Delta f/f_o}$$

Generalized reflection coefficient

$$\Gamma(\lambda_o/4) = \Gamma_L e^{-j(4\pi/\lambda)\lambda_o/4} = \Gamma_L e^{-j\pi\lambda_o/\lambda} = \Gamma_L e^{-j\pi\lambda_o/\lambda}$$

$$\Rightarrow \Gamma(\lambda_o/4) = \Gamma_L e^{-j\pi(1+\Delta f/f_o)} = -\Gamma_L e^{-j\pi\Delta f/f_o}$$

Calculate $Z(\Delta f)$ – Step 2



Impedance calculation

$$Z(\lambda_o/4) = Z_o \frac{1 - \Gamma_L e^{-j\pi\Delta f/f_o}}{1 + \Gamma_L e^{-j\pi\Delta f/f_o}} = Z_o \frac{1 + e^{-j\pi\Delta f/f_o}}{1 - e^{-j\pi\Delta f/f_o}}$$

• Recall $e^{-j\pi\Delta f/f_o} pprox 1 - j\pi\Delta f/f_o$

$$\Rightarrow Z(z) \approx Z_o \frac{1 + 1 - j\pi\Delta f/f_o}{1 - 1 + j\pi\Delta f/f_o} \approx Z_o \frac{2}{j\pi\Delta f/f_o} = -j\frac{2}{\pi}\sqrt{\frac{L}{C}} \left(\frac{w_o}{\Delta w}\right)$$

Looks like LC tank circuit (but more than one mode)!

Smith Chart

Define normalized impedance

$$Z_n = \frac{Z_L}{Z_o}$$

Mapping from normalized impedance to Γ is one-to-one

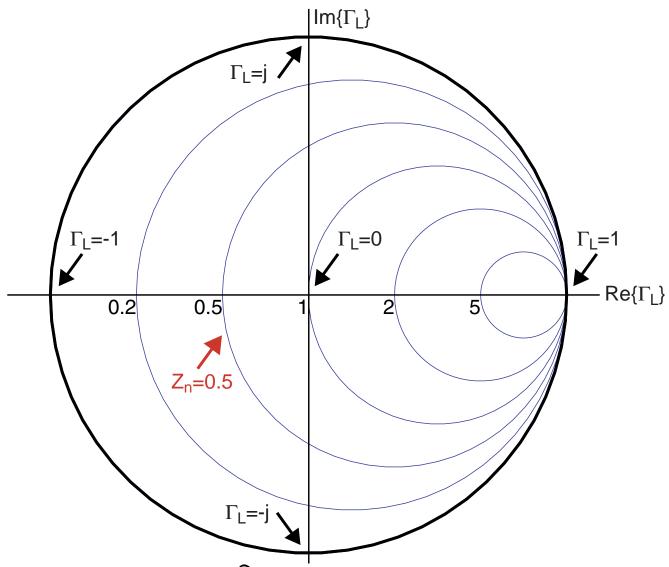
$$Z_n = \frac{1 + \Gamma_L}{1 - \Gamma_L}$$

- $lue{}$ Consider working in coordinate system based on Γ
- **Key relationship between Z_n and \Gamma**

$$Re\{Z_n\} + jIm\{Z_n\} = \frac{1 + Re\{\Gamma_L\} + jIm\{\Gamma_L\}}{1 - (Re\{\Gamma_L\} + jIm\{\Gamma_L\})}$$

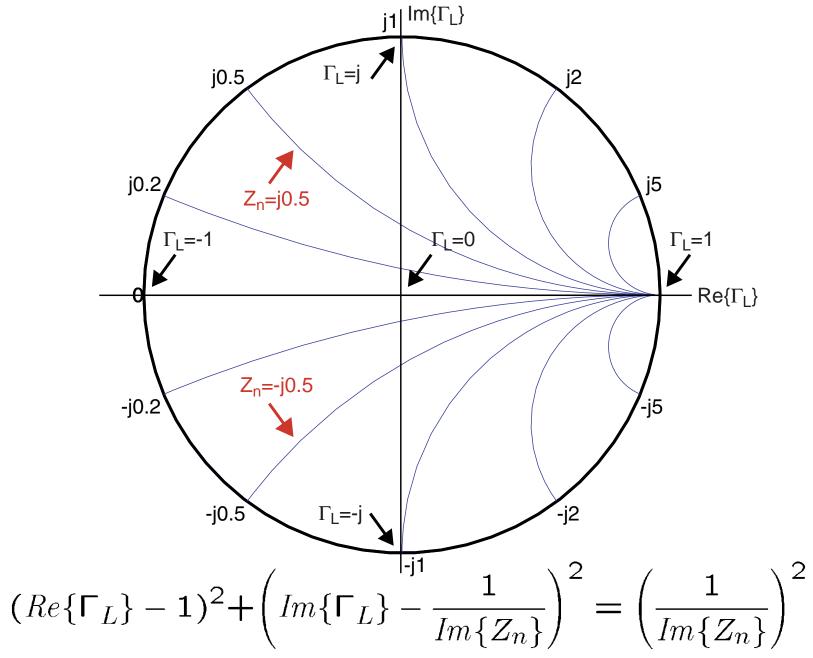
Equate real and imaginary parts to get Smith Chart

Real Impedance in \(\Gamma\) Coordinates (Equate Real Parts)



$$\left(Re\{\Gamma_L\} - \frac{Re\{Z_n\}}{1 + Re\{Z_n\}}\right)^2 + (Im\{\Gamma_L\})^2 = \left(\frac{1}{1 + Re\{Z_n\}}\right)^2$$

Imag. Impedance in \(\Gamma\) Coordinates (Equate Imag. Parts)



What Happens When We Invert the Impedance?

Fundamental formulas

$$Z_n = \frac{1 + \Gamma_L}{1 - \Gamma_L} \Rightarrow \Gamma_L = \frac{Z_n - 1}{Z_n + 1}$$

Impact of inverting the impedance

$$Z_n \to 1/Z_n \Rightarrow \Gamma_L \to -\Gamma_L$$

Derivation:

$$\frac{1/Z_n - 1}{1/Z_n + 1} = \frac{1 - Z_n}{1 + Z_n} = -\left(\frac{Z_n - 1}{Z_n + 1}\right)$$

- We can invert complex impedances in Γ plane by simply changing the sign of Γ!
- How can we best exploit this?

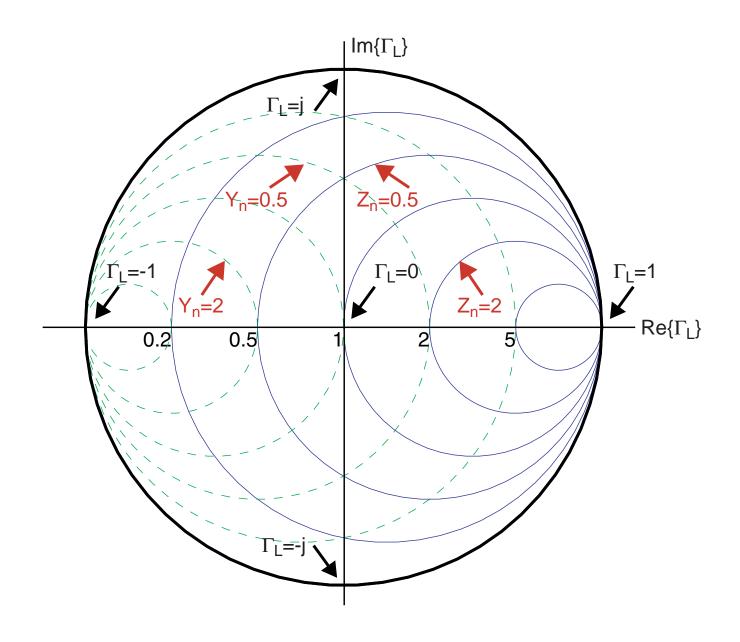
The Smith Chart as a Calculator for Matching Networks

 Consider constructing both impedance and admittance curves on Smith chart

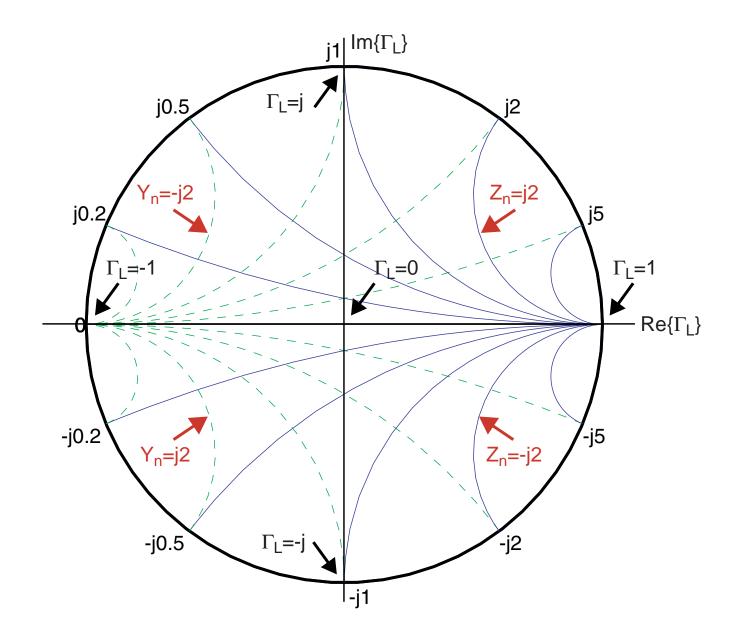
$$Z_n \to 1/Z_n \Rightarrow \Gamma_L \to -\Gamma_L$$

- Conductance curves derived from resistance curves
- Susceptance curves derived from reactance curves
- For series circuits, work with impedance
 - Impedances add for series circuits
- For parallel circuits, work with admittance
 - Admittances add for parallel circuits

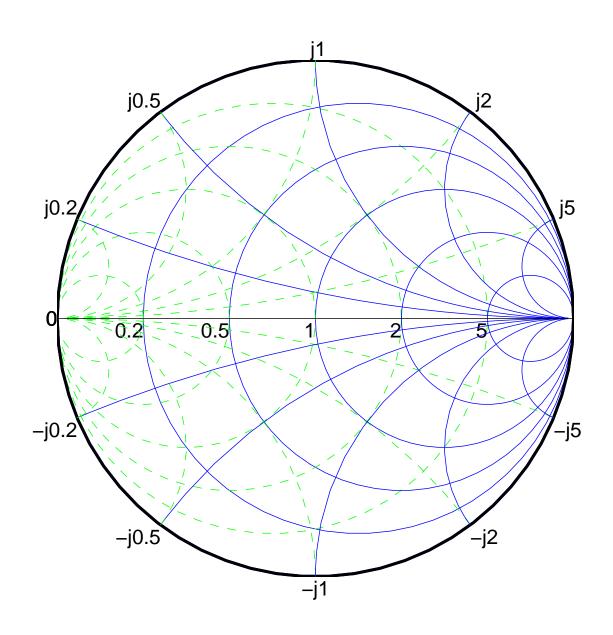
Resistance and Conductance on the Smith Chart



Reactance and Susceptance on the Smith Chart

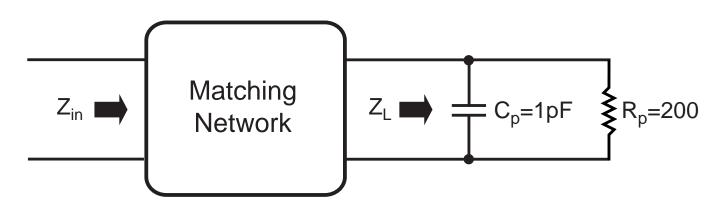


Overall Smith Chart



Example – Match RC Network to 50 Ohms at 2.5 GHz

Circuit

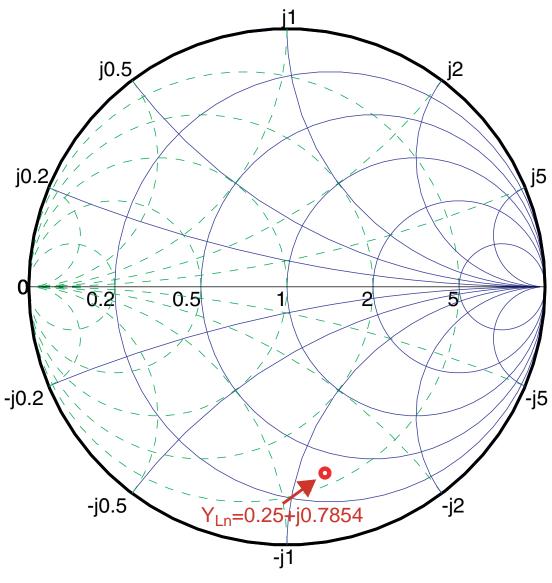


Step 1: Calculate Z_{Ln}

$$Z_{Ln} = \frac{Z_L}{Z_o} = \frac{R_L || (1/jwC)}{50} = \frac{1}{50(1/R_L + jwC)}$$
$$= \frac{1}{50(1/200 + j2\pi(2.5e9)10^{-12})} = \frac{1}{0.25 + j.7854}$$

Step 2: Plot Z_{Ln} on Smith Chart (use admittance, Y_{Ln})

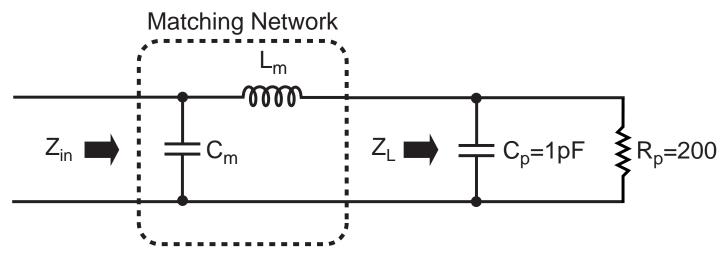
Plot Starting Impedance (Admittance) on Smith Chart



(Note: $Z_{Ln}=0.37-j1.16$)

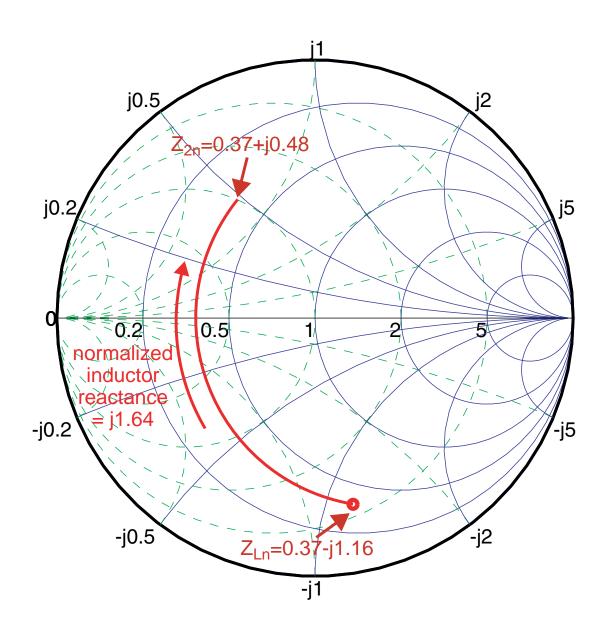
Develop Matching "Game Plan" Based on Smith Chart

By inspection, we see that the following matching network can bring us to Z_{in} = 50 Ohms (center of Smith chart)



- Use the Smith chart to come up with component values
 - Inductance L_m shifts impedance up along reactance curve
 - Capacitance C_m shifts impedance down along susceptance curve

Add Reactance of Inductor L_m



Inductor Value Calculation Using Smith Chart

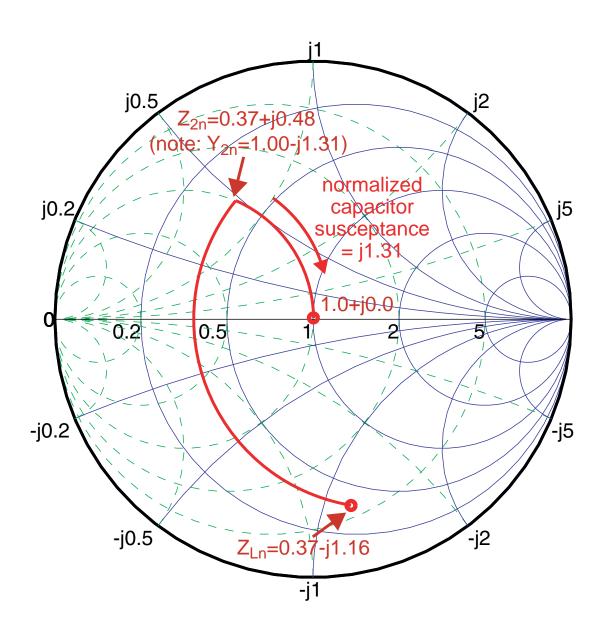
 From Smith chart, we found that the desired normalized inductor reactance is

$$\frac{jwL_m}{Z_0} = \frac{jwL_m}{50} = j1.64$$

Required inductor value is therefore

$$\Rightarrow L_m = \frac{50(1.64)}{2\pi 2.5e9} = 5.2nH$$

Add Susceptance of Capacitor C_m (Achieves Match!)



Capacitor Value Calculation Using Smith Chart

From Smith chart, we found that the desired normalized capacitor susceptance is

$$Z_{o}jwC_{m} = 50jwC_{m} = j1.31$$

Required capacitor value is therefore

$$\Rightarrow C_m = \frac{1.31}{50(2\pi 2.5e9)} = 1.67pF$$

Just For Fun

Play the "matching game" at

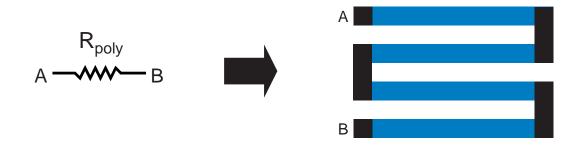
http://contact.tm.agilent.com/Agilent/tmo/an-95-1/classes/imatch.html

- Allows you to graphically tune several matching networks
- Note: game is set up to match source to load impedance rather than match the load to the source impedance
 - Same results, just different viewpoint

Passives

Polysilicon Resistors

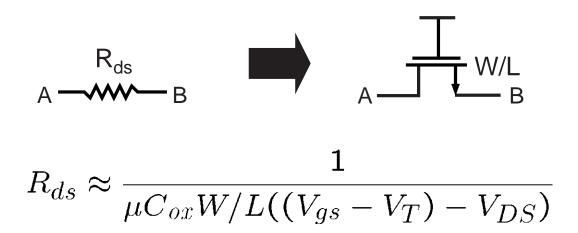
Use unsilicided polysilicon to create resistor



- Key parameters
 - Resistance (usually 100- 200 Ohms per square)
 - Parasitic capacitance (usually small)
 - Appropriate for high speed amplifiers
 - Linearity (quite linear compared to other options)
 - \blacksquare Accuracy (usually can be set within \pm 15%)

MOS Resistors

Bias a MOS device in its triode region



- High resistance values can be achieved in a small area (MegaOhms within tens of square microns)
- Resistance is quite nonlinear
 - Appropriate for small swing circuits

High Density Capacitors (Biasing, Decoupling)

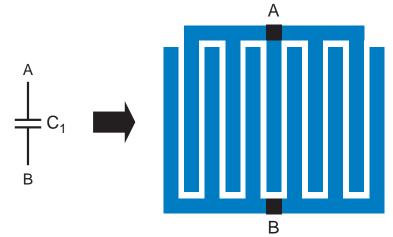
- MOS devices offer the highest capacitance per unit area
 - Limited to a one terminal device
 - Voltage must be high enough to invert the channel



- Key parameters
 - Capacitance value
 - Raw cap value from MOS device is 6.1 fF/μ m² for 0.24u
 CMOS
 - Q (i.e., amount of series resistance)
 - Maximized with minimum L (tradeoff with area efficiency)
- See pages 39-40 of Tom Lee's book

High Q Capacitors (Signal Path)

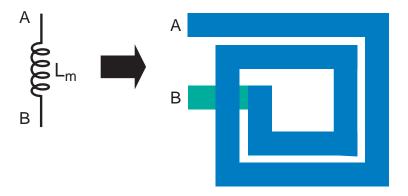
- Lateral metal capacitors offer high Q and reasonably large capacitance per unit area
 - Stack many levels of metal on top of each other (best layers are the top ones), via them at maximum density



- Accuracy often better than $\pm 10\%$
- Parasitic side cap is symmetric, less than 10% of cap value
- **Example:** $C_T = 1.5$ fF/μm² for 0.24μm process with 7 metals, $L_{min} = W_{min} = 0.24$ μm, $t_{metal} = 0.53$ μm
 - see "Capacity Limits and Matching Properties of Integrated Capacitors", Aparicio et. al., JSSC, Mar 2002

Spiral Inductors

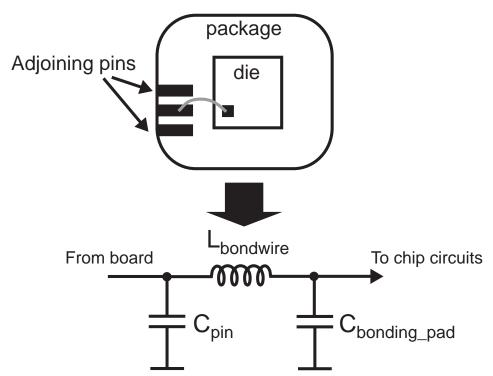
 Create integrated inductor using spiral shape on top level metals (may also want a patterned ground shield)



- Key parameters are Q (< 10), L (1-10 nH), self resonant freq.</p>
- Usually implemented in top metal layers to minimize series resistance, coupling to substrate
- Design using Mohan et. al, "Simple, Accurate Expressions for Planar Spiral Inductances, JSSC, Oct, 1999, pp 1419-1424
- Verify inductor parameters (L, Q, etc.) using ASITIC
 http://formosa.eecs.berkeley.edu/~niknejad/asitic.html

Bondwire Inductors

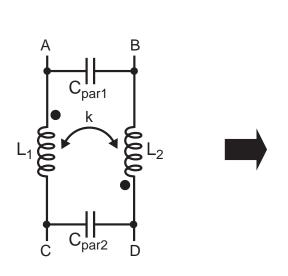
- Used to bond from the package to die
 - Can be used to advantage

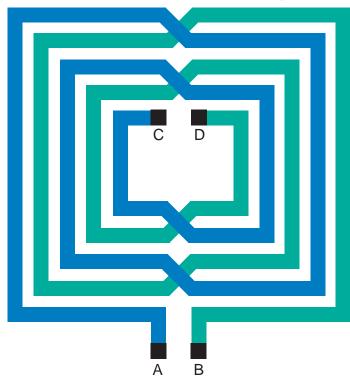


- Key parameters
 - Inductance (\approx 1 nH/mm usually achieve 1-5 nH)
 - Q (much higher than spiral inductors typically > 40)

Integrated Transformers

Utilize magnetic coupling between adjoining wires





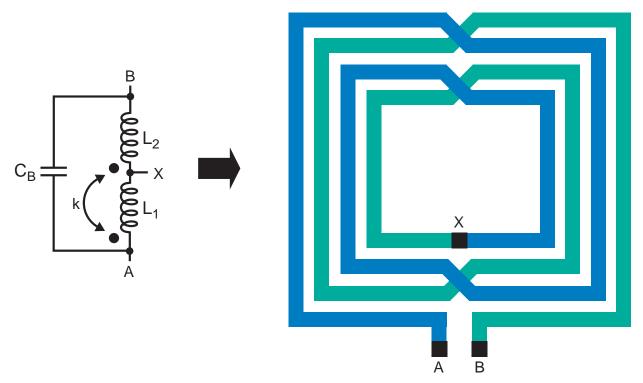
- Key parameters
 - L (self inductance for primary and secondary windings)
 - k (coupling coefficient between primary and secondary)

Note: $k = \frac{M}{\sqrt{L_1 L_2}}$ where M = mutual inductance

Design – ASITIC, other CAD packages

High Speed Transformer Example – A T-Coil Network

 A T-coil consists of a center-tapped inductor with mutual coupling between each inductor half



- Used for bandwidth enhancement
 - See S. Galal, B. Ravazi, "10 Gb/s Limiting Amplifier and Laser/Modulator Driver in 0.18u CMOS", ISSCC 2003, pp 188-189 and "Broadband ESD Protection ...", pp. 182-183